

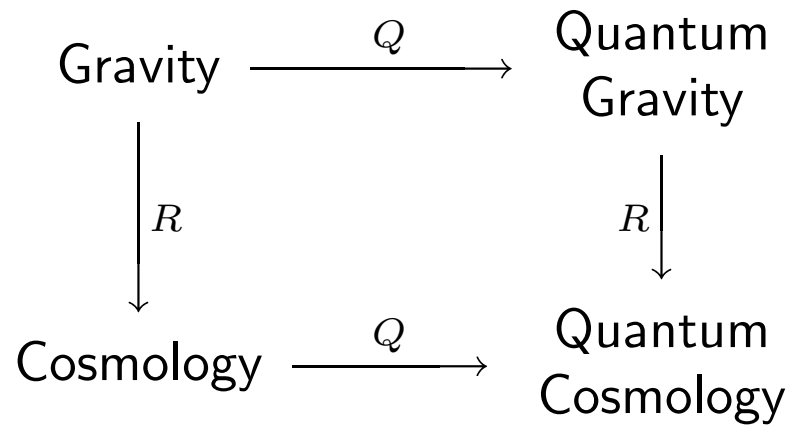
On the Configuration Spaces of Homogeneous Loop Quantum Cosmology and Loop Quantum Gravity

Christian Fleischhack

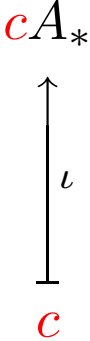
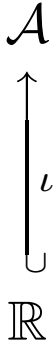
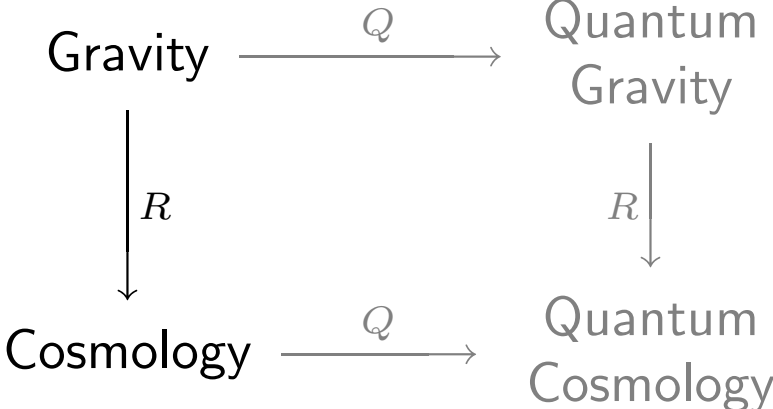
Department Mathematik
Universität Hamburg

Zakopane, March 2008

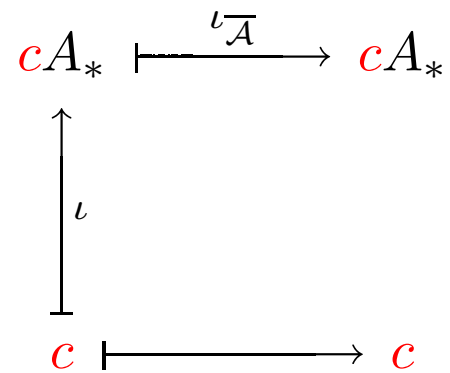
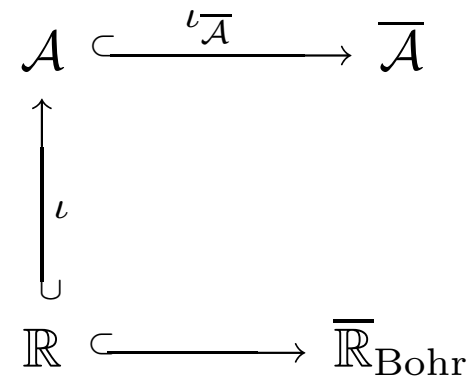
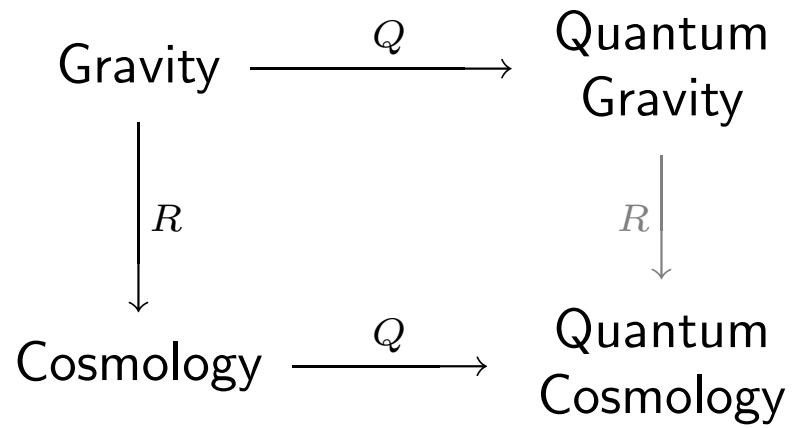
1 Motivation



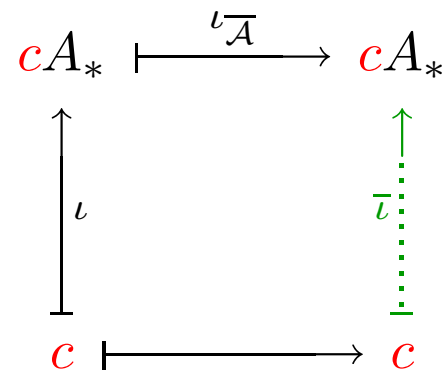
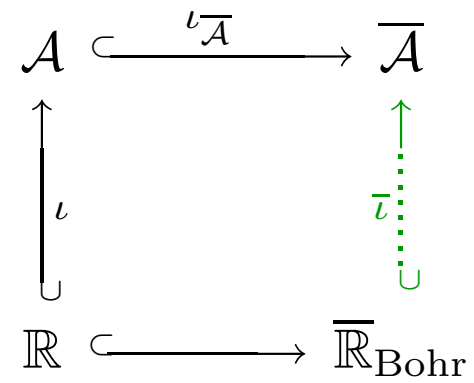
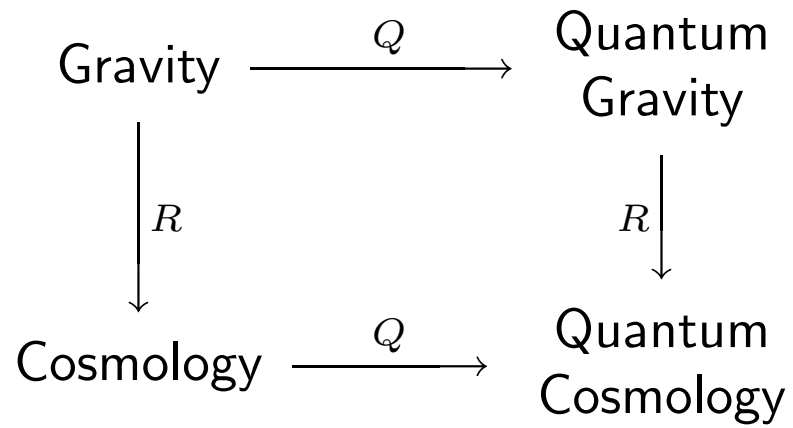
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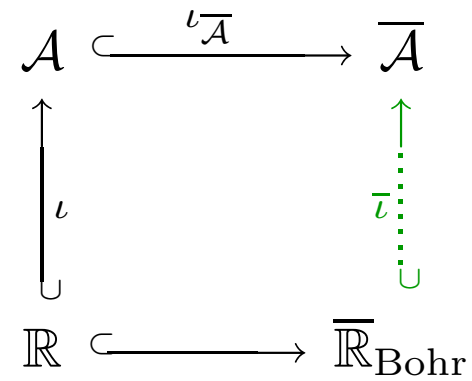
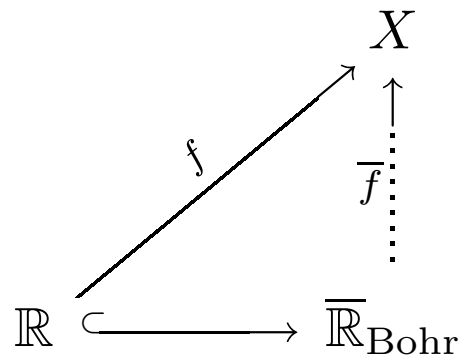
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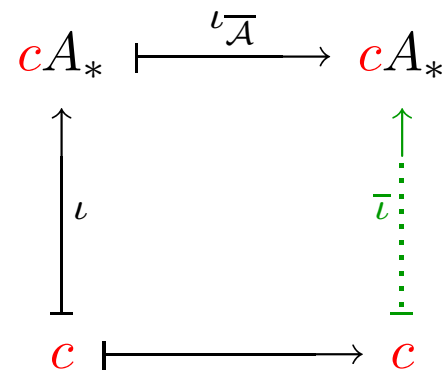


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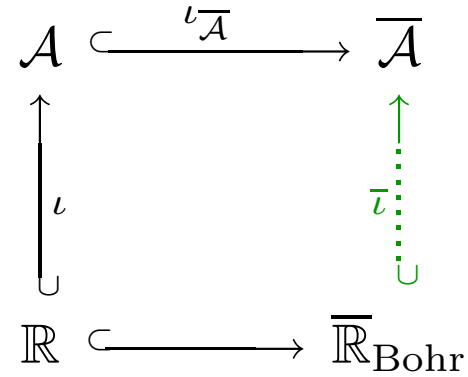
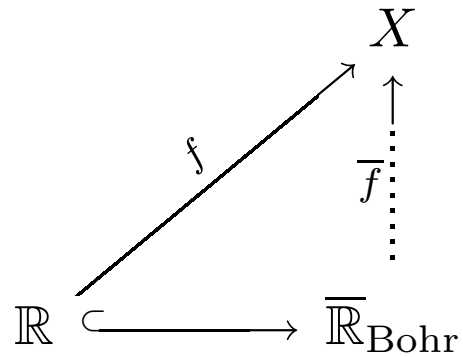


Theorem:

f continuous w.r.t. Bohr topology
 $\iff f$ almost periodic



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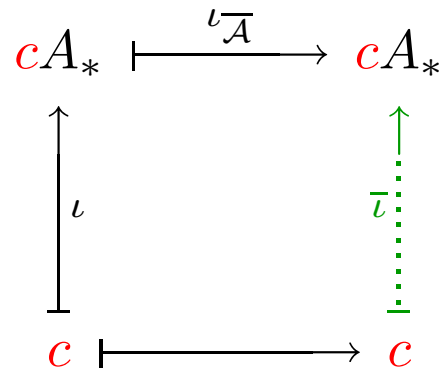


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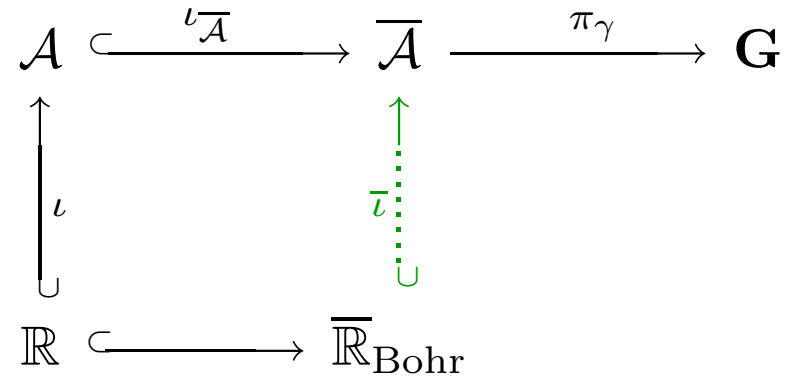
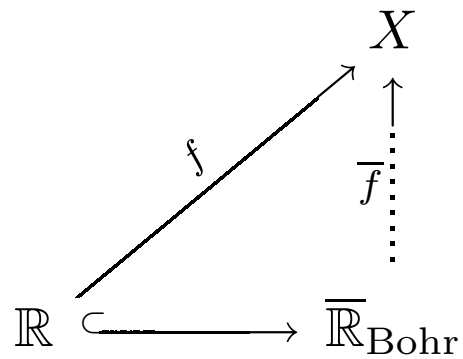
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Corollary:

\bar{l} exists
 $\iff \bar{l} \circ \iota_{\bar{A}}$ continuous



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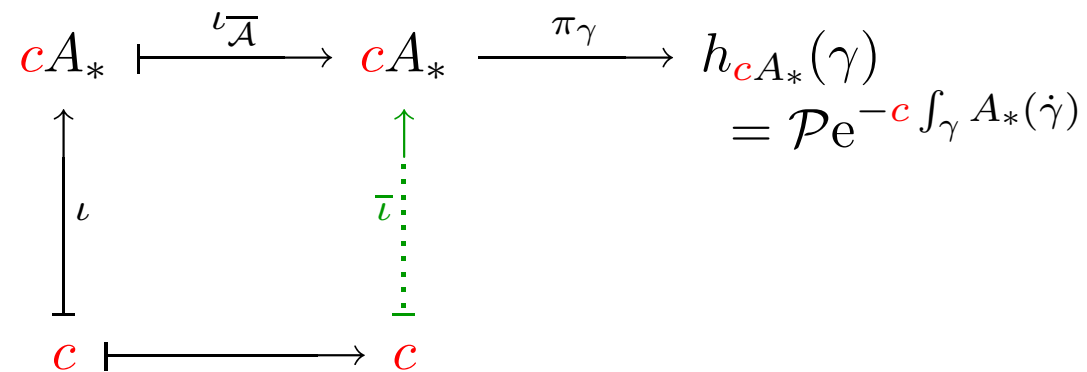
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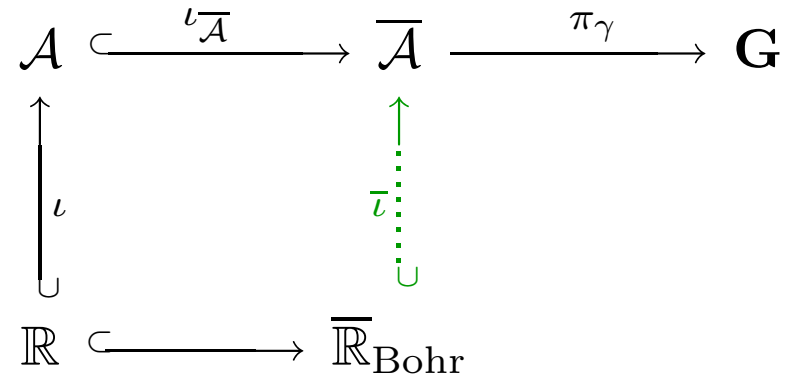
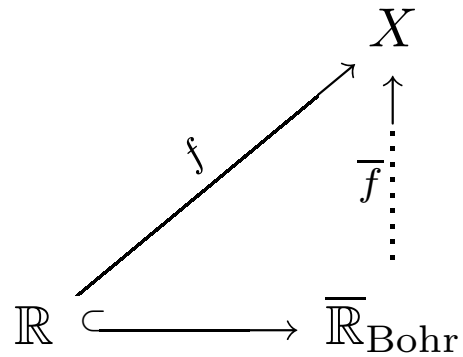
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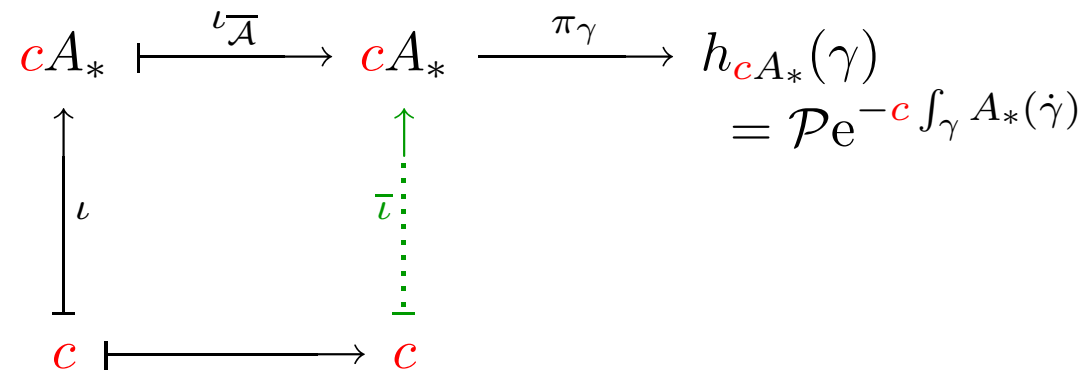
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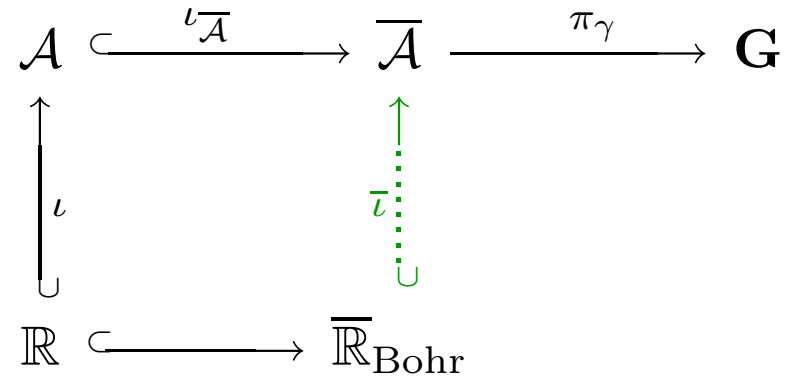
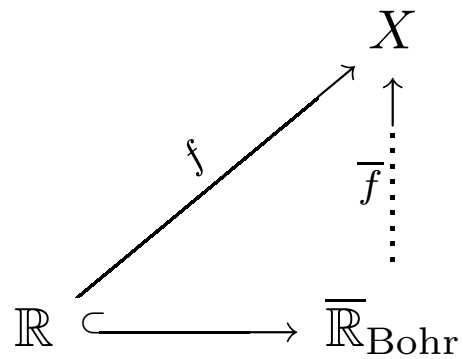
Corollary:

$\bar{\iota}$ exists

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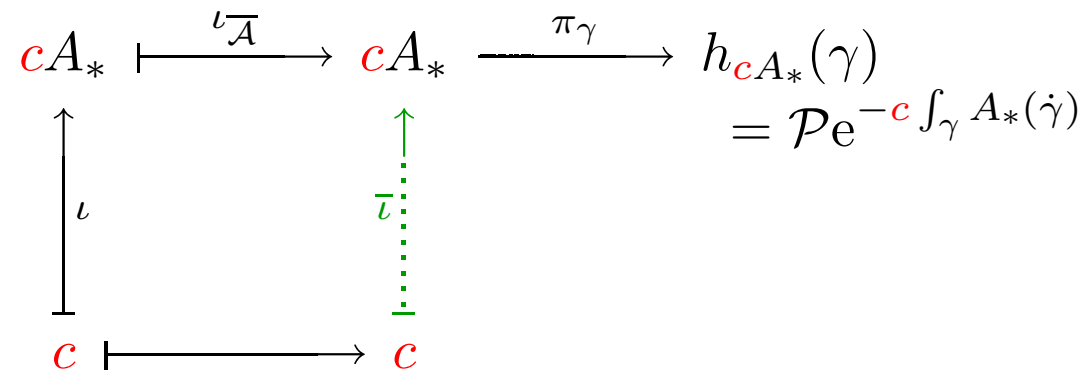
Corollary:

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Question: For which γ is $c \mapsto h_{cA_*}(\gamma)$ almost periodic?

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- Parallel Transport

$$g(t) := h_{cA_*}(\gamma|_{[0,t]})$$



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$$\begin{aligned}\dot{g}(t) &= -c A_*(\dot{\gamma}(t)) g(t) \\ g(0) &= \mathbf{1}\end{aligned}$$

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$$(M = \mathbb{R}^3, \mathbf{G} = SU(2))$$

$$A_* = \tau_1 dx + \tau_2 dy + \tau_3 dz$$

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$$A_*(\dot{\gamma}(t)) = \begin{pmatrix} -i\dot{z} & -i\dot{x} - \dot{y} \\ -i\dot{x} + \dot{y} & i\dot{z} \end{pmatrix}$$

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$$A_*(\dot{\gamma}(t)) = -i \begin{pmatrix} n & m \\ \bar{m} & -n \end{pmatrix}$$

with

$$\begin{aligned} m &:= \dot{x} - i\dot{y} \\ n &:= \dot{z} \end{aligned}$$

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- Differential Equation

$$\begin{pmatrix} \dot{a} & \dot{b} \\ -\dot{b} & \dot{a} \end{pmatrix} = ic \begin{pmatrix} n & m \\ \bar{m} & -n \end{pmatrix} \begin{pmatrix} a & b \\ -\bar{b} & \bar{a} \end{pmatrix}$$

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$$\begin{aligned} \dot{a} &= ic(na - m\bar{b}) & a(0) &= 1 \\ \dot{b} &= ic(nb + m\bar{a}) & b(0) &= 0 \end{aligned}$$

3 Examples

Question: For which γ is $c \mapsto h_{cA_*}(\gamma)$ almost periodic?

- Second-order Equation

$$\begin{array}{l} \dot{a} = ic(na - m\bar{b}) \\ \dot{b} = ic(nb + m\bar{a}) \end{array}$$

$$\ddot{a} = ic(\dot{n} - Mn)a - c^2a + M\dot{a}$$

with $M := \frac{\dot{m}}{m}$

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n, \dot{n}		
Equations		
Initial Values ($m(0) = -i$)		
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Equations	$\ddot{a} + c^2a = 0$ $\ddot{b} + c^2b = 0$	
Initial Values ($m(0) = -i$)	$b(0) = 0$ $\dot{b}(0) = c$	
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	Straight Line	Planar Circle
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Initial Values ($m(0) = -i$)	$b(0) = 0$ $\dot{b}(0) = c$	
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Theorem: $\iota : \mathbb{R} \longrightarrow \mathcal{A}$ cannot be continuously extended to $\bar{\iota} : \overline{\mathbb{R}}_{\text{Bohr}} \longrightarrow \overline{\mathcal{A}}$.

4 Non-Almost Periodicity of Parallel Transports

Question: For which γ is $c \mapsto h_{cA_*}(\gamma)$ almost periodic?

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Question: For which γ is $c \mapsto h_{cA_*}(\gamma)$ almost periodic?

Lemma: We have $\lim_{c \rightarrow \infty} (\bar{m}a^2 + 2nab\bar{b} - m\bar{b}^2) \Big|_0^t = 0$.

$$\begin{array}{l} \dot{a} = ic(na - m\bar{b}) \\ \dot{b} = ic(nb + m\bar{a}) \end{array}$$

Definition: t almost periodic $\iff a(t)$ and $b(t)$ almost periodic.

Corollary: t almost periodic $\implies F(t) := (\bar{m}a^2 + 2nab\bar{b} - m\bar{b}^2) \Big|_0^t = 0$

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Then γ is a straight line.

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Proposition: Let there exist an accumulation point of almost periodic t .
Then γ is a straight line.

Proof: Assume t almost periodic with $\dot{m}(t) \neq 0$. Then $\dot{m} \neq 0$ on some open $U \ni t$.

$$\begin{aligned} \text{Assumption} \implies F \equiv 0 &\implies \dot{\bar{m}}a^2 + 2\dot{n}a\bar{b} - \dot{m}\bar{b}^2 = \dot{F} \equiv 0 \\ &\qquad \qquad \qquad \bar{m}a^2 + 2nab - m\bar{b}^2 = m(0) \end{aligned}$$

$$\implies a \text{ and } b \text{ independent of } c \text{ on } U; \text{ as well as } \dot{a} \text{ and } \dot{b}$$

$$\implies na - m\bar{b} = 0 = nb + m\bar{a}$$

$$\implies a = (n^2 + \bar{m}m)a = nm\bar{b} + \bar{m}ma = 0 \quad \zeta$$

Thus $\dot{m}(t) = 0$ for all almost periodic $t \implies \dot{m} \equiv 0 \implies \dot{n} \equiv 0$. **qed**

4 Non-Almost Periodicity of Parallel Transports

Question: For which γ is $c \mapsto h_{cA_*}(\gamma)$ almost periodic?

Lemma: We have $\lim_{c \rightarrow \infty} (\bar{m}a^2 + 2nab\bar{b} - m\bar{b}^2) \Big|_0^t = 0$.

$$\begin{aligned} \dot{a} &= ic(na - m\bar{b}) \\ \dot{b} &= ic(nb + m\bar{a}) \end{aligned}$$

Proof: Recall $\ddot{a} = ic(\dot{n} - Mn)a - c^2a + M\dot{a}$, and observe, with $\bar{m}m + n^2 = 1$,

$$\bar{m}a^2 + 2nab\bar{b} - m\bar{b}^2 = \frac{1}{m} \left(-(na - m\bar{b})^2 + a^2 \right) = \frac{1}{m} \left(\frac{\dot{a}^2}{c^2} + a^2 \right).$$

Use

$$\begin{aligned} a =: \sqrt{m} \alpha &\implies \ddot{\alpha} + c^2\alpha = \left(\frac{1}{4}M^2 - \frac{1}{2}\dot{M} + ic(\dot{n} - Mn) \right) \alpha =: \rho\alpha \\ &\implies \frac{d}{dt} \left(\frac{\dot{\alpha}^2}{c^2} + \alpha^2 \right) = \frac{2}{c^2} (\ddot{\alpha} + c^2\alpha) \dot{\alpha} = \frac{\rho}{c^2} \frac{d}{dt} \alpha^2. \end{aligned}$$

Now,

$$\left(\frac{\dot{\alpha}^2}{c^2} + \alpha^2 \right) \Big|_0^t = - \int_0^t \frac{\dot{\rho}}{c^2} \alpha^2 d\tau + \frac{\rho}{c^2} \alpha^2 \Big|_0^t \xrightarrow{c \rightarrow \infty} 0$$

and

$$\frac{1}{m} \left(\frac{\dot{a}^2}{c^2} + a^2 \right) \Big|_0^t = \left(\frac{\dot{\alpha}^2}{c^2} + \alpha^2 \right) \Big|_0^t + \frac{Ma}{mc^2} \left(\dot{a} - \frac{1}{4}Ma \right) \Big|_0^t \xrightarrow{c \rightarrow \infty} 0.$$

qed

5 Conclusions

Brunnemann, CF: 0709.1621

- **Theorem:** Let γ be an analytic curve that is not part of a straight line. Then there is some $T > 0$, such that the parallel transport along $\gamma|_{[0,t]}$ is **not** almost periodic w.r.t. c for any $0 < t < T$.

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- **Corollary:** The configuration space of (**homogeneous, isotropic**) LQC **cannot** be continuously embedded into that of LQG by extension of the classical embedding:

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- **Open Questions:**
 - almost periodicity “by chance”?
 - asymptotic almost periodicity
 - $k = \pm 1$
 - general Lie groups
 - ...